

Solving PDEs on Overlapping Grids with Overture

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Major Contributors

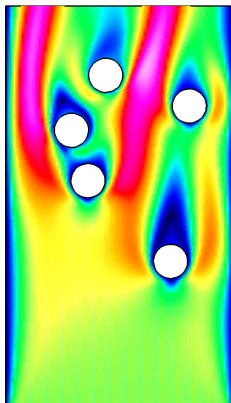
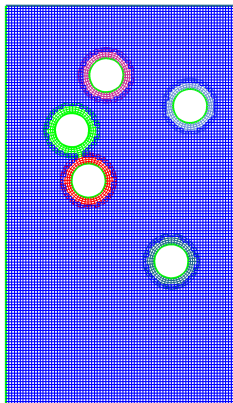
Don Schwendeman (RPI),

Jeff Banks (LLNL).



What are overlapping grids and why are they useful?

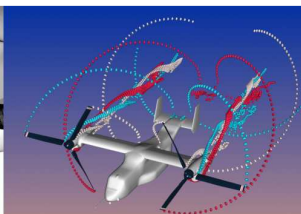
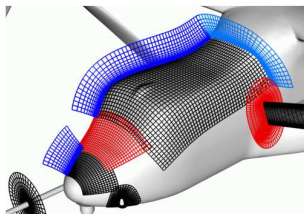
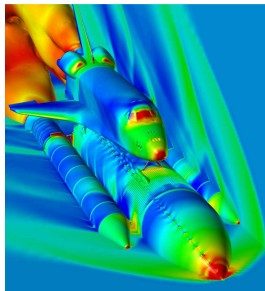
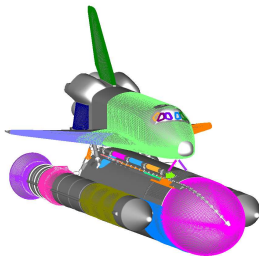
Overlapping grid: a set of structured grids that overlap.



- Overlapping grids can be rapidly generated as bodies move.
- High quality grids under large displacements.
- Cartesian grids for efficiency.
- Efficient for high-order accurate methods.



Aerospace applications using overlapping grids.



Space shuttle figures courtesy of William Chan and Reynaldo Gomez.

V-22 Osprey figures courtesy of William Chan, Andrew Wissink and Robert Meakin.



Overture: tools for solving PDE's on overlapping grids

- high level C++ interface for rapid prototyping of PDE solvers.
- built upon optimized C and fortran kernels.
- library of finite-difference operators: conservative and non-conservative, 2nd, 4th, 6th and 8th order accurate approximations.
- support for moving grids.
- support for block structured adaptive mesh refinement (AMR).
- extensive grid generation capabilities (Ogen).
- CAD fixup tools (for CAD from IGES files).
- interactive graphics and data base support (HDF).



The CG (Composite Grid) suite of PDE solvers

- **cgad**: advection diffusion equations.
- **cgins**: incompressible Navier-Stokes with heat transfer.
- **cgcns**: compressible Navier-Stokes, reactive Euler equations.
- **cgmp**: multi-physics solver (e.g. conjugate heat transfer).
- **cgmx**: time domain Maxwell's equations solver.
- **cgsm**: solid mechanics (*new*)

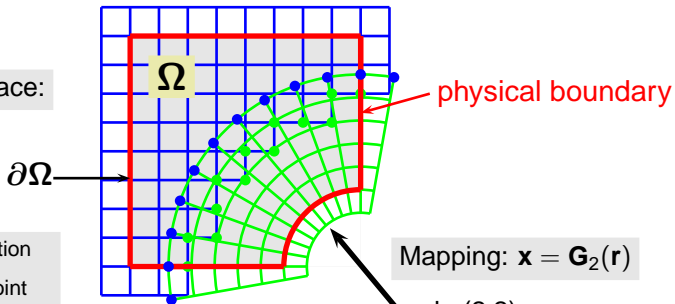
Overture and CG are freely available from the web:

www.llnl.gov/CASC/Overture



Components of an Overlapping Grid

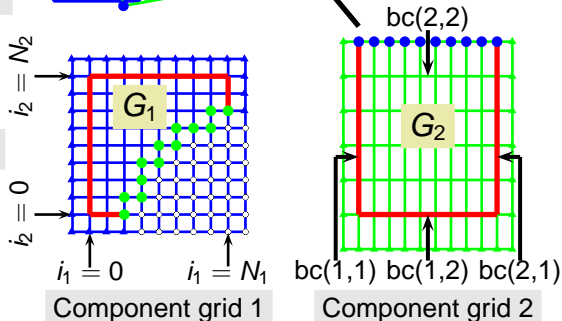
Physical space:



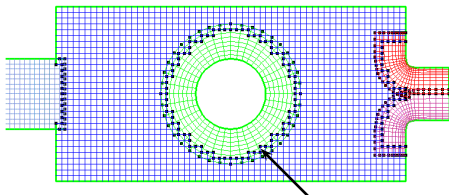
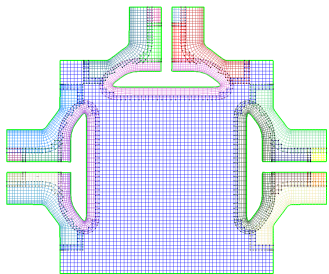
- interpolation
- unused
- ▲ ghost point

Mapping: $\mathbf{x} = \mathbf{G}_2(\mathbf{r})$

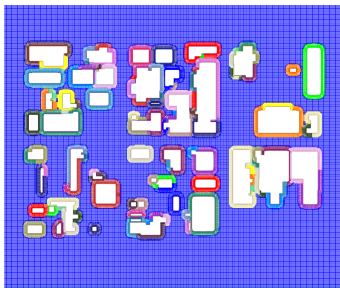
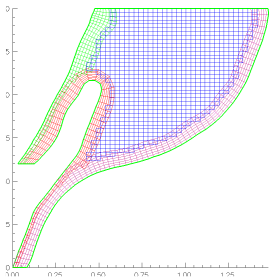
Parameter space:



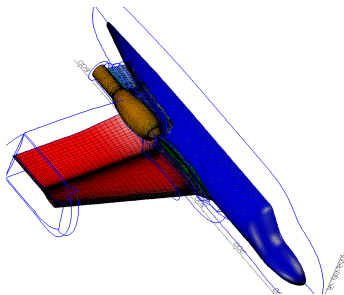
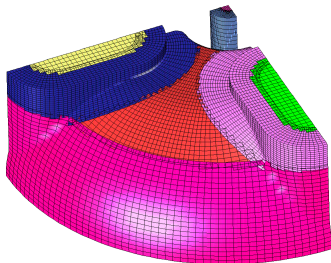
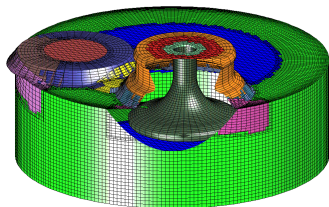
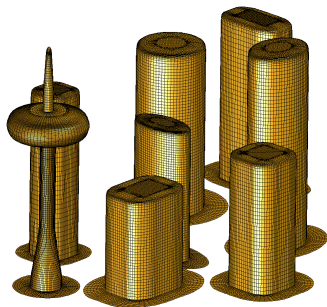
Ogen can be used to build 2D overlapping grids:



Solutions coupled by interpolation



Ogen can be used to build 3D overlapping grids:



Composite/ Chimera/ Overset/ Overlapping Grids

A Short History

- Volkov, circa [1966] developed a *Composite Mesh* method for Laplace's equation on regions with piece-wise smooth boundaries separated by corners. Polar grids are fitted at corners to handle potential singularities.
- Starius, circa [1977] (student of H.-O. Kreiss) considered *Composite Mesh* methods for elliptic and hyperbolic problems – introduces a hyperbolic grid generator.
- Steger, circa [1980] independently conceives the idea of the overlapping grid, subsequently named the *Chimera* approach after the mythical Chimera beast having a human face, a lion's mane and legs, a goat's body, and dragon's tail. NASA groups develop grid generator PEGSUS, hyperbolic grid generation and flow solver Overflow (Steger, Benek, Suhs, Buning, Chan, Meakin, et. al.)
- B. Kreiss circa [1980] develops overlapping grid generator which subsequently leads to the CMPGRD grid generator [1983] (Chesshire, Henshaw) later leading to the Overture set of tools [1994].



Theory for finite difference schemes

There is extensive numerical analysis theory underpinning this work.

- classic von Neumann stability analysis (periodic domains).
- energy estimates (L_2 -norm estimates).
- normal mode analysis, GKS theory (initial boundary value problems).

Some references:

- Gustafsson, Kreiss, Olinger, *Time Dependent Problems and Difference Methods*, (book).
- Strikwerda, *Finite Difference Schemes and Partial Differential Equations*, (book).
- Gustafsson, Kreiss, Sundström, *Stability Theory of Difference Approx. for Mixed Initial Boundary Value Problems, I. and II.*, Math. Comp.
- Starius, *On Composite Mesh Difference Methods for Hyperbolic Differential Equations*, Numer. Math.



A one-dimensional overlapping grid example

To solve the advection-diffusion equation

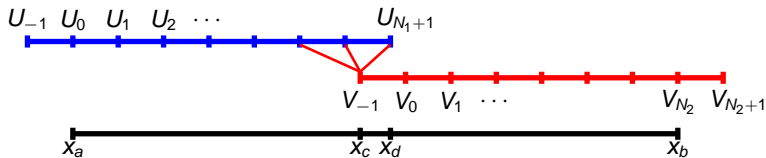
$$\begin{aligned}u_t + au_x &= \nu u_{xx}, & x \in (0, 1) \\u(0, t) &= g_0(t), \quad u_x(1, t) = g_1(t), & \text{(boundary conditions)} \\u(x, 0) &= u_0(x), & \text{(initial conditions)}\end{aligned}$$

introduce grid points on the two overlapping component grids,

$$x_i^{(1)} = x_a + i\Delta x_1, \quad i = -1, 0, 1, \dots, N_1 + 1, \quad \Delta x_1 = (x_d - x_a)/N_1$$

$$x_j^{(2)} = x_c + (j + 1)\Delta x_2, \quad j = -1, 0, 1, \dots, N_2 + 1, \quad \Delta x_2 = (x_b - x_c)/N_2$$

and approximations $U_i^n \approx u(x_i^{(1)}, n\Delta t)$, $V_j^n \approx u(x_j^{(2)}, n\Delta t)$.



Discretize with forward-Euler and central differences

Given the solution at time t^n , compute the solution at time t^{n+1} :

$$(U_i^{n+1} - U_i^n)/\Delta t = -aD_0U_i^n + \nu D_+D_-U_i^n, \quad i = 1, 2, \dots, N_1$$

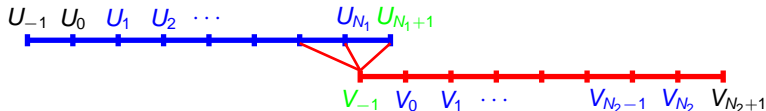
$$(V_j^{n+1} - V_j^n)/\Delta t = -aD_0V_j^n + \nu D_+D_-V_j^n, \quad j = 0, 2, \dots, N_2$$

$$U_0^{n+1} = g(t^n), \quad D_0V_{N_2}^{n+1} = g_1(t^{n+1}), \quad (\text{boundary conditions})$$

$$U_{N_1+1}^{n+1} = (1 - \alpha)(1 - \frac{\alpha}{2}) V_{-1}^{n+1} + \alpha(2 - \alpha) V_0^{n+1} + \frac{\alpha}{2}(\alpha - 1) V_1^{n+1}, \quad (\text{interpolation})$$

$$V_{-1}^{n+1} = (1 - \beta)(1 - \frac{\beta}{2}) U_{N_1-1}^{n+1} + \beta(2 - \beta) U_{N_1}^{n+1} + \frac{\beta}{2}(\beta - 1) U_{N_1+1}^{n+1}, \quad (\text{interpolation})$$

$$D_0U_i^n = \frac{U_{i+1}^n - U_{i-1}^n}{2\Delta x}, \quad D_+U_i^n = \frac{U_{i+1}^n - U_i^n}{\Delta x}, \quad D_-U_i^n = \frac{U_i^n - U_{i-1}^n}{\Delta x}.$$



Overture supports a high-level C++ interface

But is built upon mainly Fortran kernels.

Solve $u_t + au_x + bu_y = \nu(u_{xx} + u_{yy})$

```
CompositeGrid cg; // create a composite grid
getFromADatabaseFile(cg,"myGrid.hdf");
floatCompositeGridFunction u(cg); // create a grid function
u=1.;
CompositeGridOperators op(cg); // operators
u.setOperators(op);
float t=0, dt=.005, a=1., b=1., nu=.1;
for( int step=0; step<100; step++ )
{
    u+=dt*( -a*u.x()-b*u.y()+nu*(u.xx()+u.yy()) ); // forward Euler
    t+=dt;
    u.interpolate();
    u.applyBoundaryCondition(0,dirichlet,allBoundaries,0.);
    u.finishBoundaryConditions();
}
```

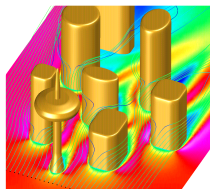
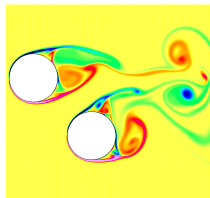


Overture is used by research groups worldwide

- Blood flow in veins with blood clot filters. (Mike Singer, LLNL).
- Pitching airfoils and micro-air vehicles (Yongsheng Lian, U. of Louisville)
- Relativistic hydrodynamics and Einstein field equations (Philip Blakely, Nikos Nikiforakis, U. Cambridge).
- Compressible flow/ice-formation (Graeme Leese, U. Cambridge).
- Tear films and droplets (Rich Braun U. Delaware, Kara Maki UMN).
- High-order accurate subsonic/transonic aero-acoustics (Phillipe Lafon, CNRS, EDF, France).
- Low Reynolds flow for pitching airfoils (D. Chandar, R. Yapalparvi, M. Damodaran, NTU, Singapore).
- Incompressible flow in pumps (J.P. Potanza, Shell Oil, Houston).
- High-order accurate, compact Hermite-Taylor schemes (Tom Hagstrom, SMU, Dallas).



Cgins: incompressible Navier-Stokes solver.



- 2nd-order and 4th-order accurate (DNS).
- support for moving rigid-bodies (not parallel yet).
- heat transfer (Boussinesq approximation).
- semi-implicit (time accurate), pseudo steady-state (efficient line solver), full implicit.

• WDH., *A Fourth-Order Accurate Method for the Incompressible Navier-Stokes Equations on Overlapping Grids*, J. Comput. Phys, **113**, no. 1, (1994) 13–25.



Incompressible Navier-Stokes.

Split-step, velocity-pressure formulation:

$$\begin{aligned} \mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p - \nu \Delta \mathbf{u} - \mathbf{f} &= 0, & t > 0, & \mathbf{x} \in \Omega \\ \nabla \cdot \mathbf{u} = 0 & \Delta p - \nabla \mathbf{u} : \nabla \mathbf{u} - \alpha \nabla \cdot \mathbf{u} - \nabla \cdot \mathbf{f} &= 0, & t > 0, \mathbf{x} \in \Omega \end{aligned}$$

Divergence damping term: $\alpha \nabla \cdot \mathbf{u}$ is important.

Wall boundary conditions:

$$\mathbf{u} = 0, \quad \nabla \cdot \mathbf{u} = 0, \quad (\text{pressure BC}) \quad \mathbf{x} \in \partial\Omega,$$

with *numerical boundary condition*:

$$p_n = -\mathbf{n} \cdot (\nu \nabla \times \nabla \times \mathbf{u}).$$

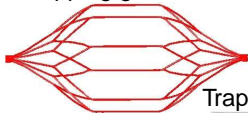
Use $\nabla \times \nabla \times \mathbf{u}$ instead of $\Delta \mathbf{u}$ for implicit time-stepping.

- WDH, N.A. Petersson, *A Split-Step Scheme for the Incompressible Navier-Stokes Equations*, 2003.

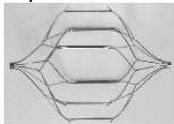


Flow past a blood-clot filter using cgins

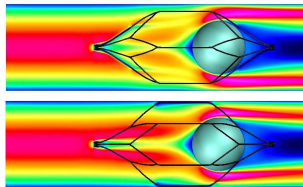
Overlapping grid for the filter



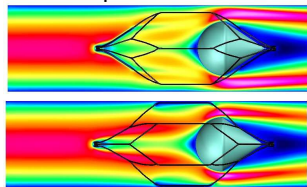
Trap-ease wire filter



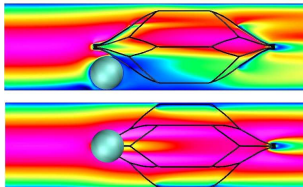
Spherical clot trapped in the filter



Cone shaped clot



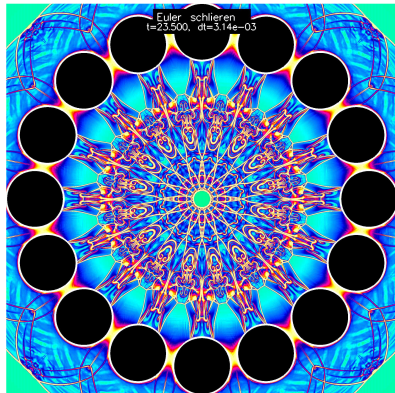
Spherical clot trapped near the front



M.A. Singer, WDH, S.L. Wang, *Computational Modeling of Blood Flow in the Trapease Inferior Vena Cava Filter*, Journal of Vascular and Interventional Radiology, **20**, 2009.



Cgcns: compressible N-S and reactive-Euler.



- reactive and non-reactive Euler equations, Don Schwendeman (RPI).
- compressible Navier-Stokes.
- multi-fluid formulation, Jeff Banks (LLNL).
- adaptive mesh refinement and moving grids.

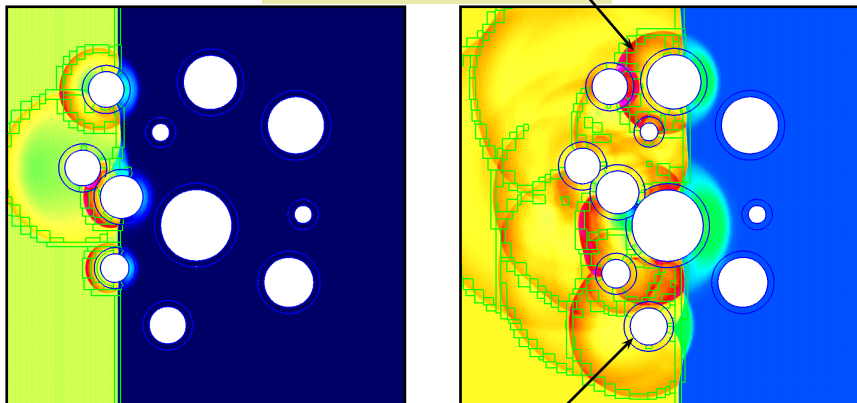
- WDH., D. W. Schwendeman, *Parallel Computation of Three-Dimensional Flows using Overlapping Grids with Adaptive Mesh Refinement*, J. Comp. Phys. **227** (2008).
- WDH., DWS, *Moving Overlapping Grids with Adaptive Mesh Refinement for High-Speed Reactive and Nonreactive Flow*, J. Comp. Phys. **216** (2005).
- WDH., DWS, *An adaptive numerical scheme for high-speed reactive flow on overlapping grids*, J. Comp. Phys. **191** (2003).



Moving overlapping grids and AMR

A shock hitting a collection of cylinders (density).

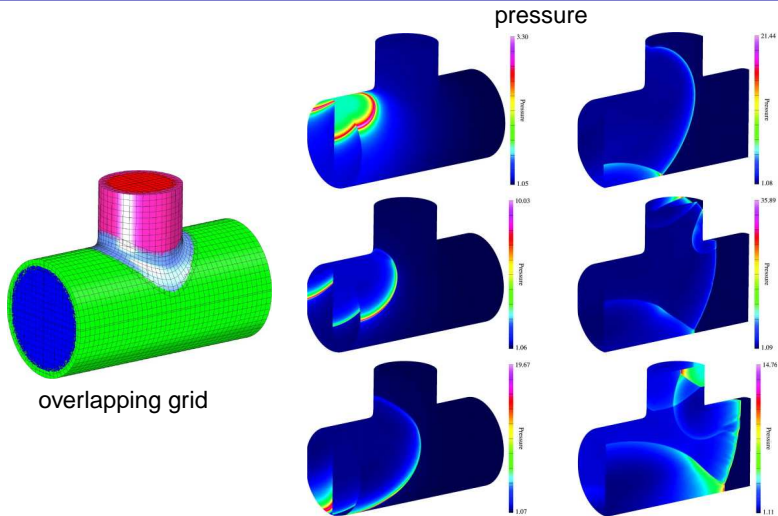
adaptive mesh refinement



moving grids



Detonation initiation in a T-shaped pipe



Notes: cgcns, reactive-Euler: one refinement level, factor 4, 4930 time steps, 48 processors, from 5 to 682 grids, 100M pts (max) (eff. resolution 400 M).



Estimating Convergence Rates

Define the volume-weighted discrete L_p -norm of a grid function U_i as

$$\|U_i\|_p = \left(\frac{\sum_i |U_i|^p dV_i}{\sum_i dV_i} \right)^{1/p}, \quad dV_i = \left| \frac{\partial \mathbf{x}}{\partial \mathbf{r}} \right|_i dr_1 dr_2 dr_3.$$

Assume the discrete solution U_i^m at grid spacing h_m satisfies

$$U_i^m - u(\mathbf{x}_i^m, t) \approx C_i^m h_m^\mu,$$

The difference between resolution h_n and h_m is

$$\|U_i^m - \mathcal{R}_n^m U_i^n\|_p \approx C |h_m^\mu - h_n^\mu|,$$

where \mathcal{R}_n^m is a fine to coarse restriction operator.

Result: Given three solutions we can estimate the convergence rate μ and the error.



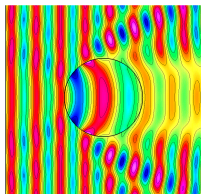
Estimating Convergence Rates

| Detonation in a T-Pipe | | | | |
|------------------------|-------------------|-------------------|-------------------|-------------------|
| | $t = 2.0$ | | $t = 2.8$ | |
| h_m | \mathcal{E}_1^m | \mathcal{E}_2^m | \mathcal{E}_1^m | \mathcal{E}_2^m |
| 1/120 | 4.0e-3 | 3.0e-2 | 3.8e-2 | 2.6e-1 |
| 1/160 | 2.2e-3 | 1.6e-2 | 2.4e-2 | 1.9e-1 |
| 1/240 | 9.8e-4 | 7.1e-3 | 1.2e-2 | 1.2e-1 |
| rate, μ | 2.04 | 2.07 | 1.65 | 1.09 |

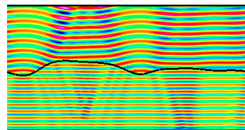
Estimated L_1 and L_2 errors in the density, \mathcal{E}_1^m and \mathcal{E}_2^m , respectively, and convergence rates μ at $t = 2.0$ and $t = 2.8$.



Cgmx: electromagnetics solver.



- a time-domain finite difference scheme.
- fourth-order accurate, 2D, 3D.
- Efficient time-stepping with the modified-equation approach
- High-order accurate symmetric difference approximations.
- High-order-accurate *centered* boundary and interface conditions.



- WDH., *A High-Order Accurate Parallel Solver for Maxwell's Equations on Overlapping Grids*, SIAM J. Scientific Computing, **28**, no. 5, (2006).



Maxwell's equations are solved in second-order form

Maxwell's equations:

$$\epsilon\mu \partial_t^2 \mathbf{E} = \Delta \mathbf{E} + \nabla (\nabla \ln \epsilon \cdot \mathbf{E}) + \nabla \ln \mu \times (\nabla \times \mathbf{E}) - \mu \partial_t \mathbf{j}$$

$$\epsilon\mu \partial_t^2 \mathbf{H} = \Delta \mathbf{H} + \nabla (\nabla \ln \mu \cdot \mathbf{H}) + \nabla \ln \epsilon \times (\nabla \times \mathbf{H}) + \epsilon \nabla \times \left(\frac{1}{\epsilon} \mathbf{j} \right)$$

Advantages of the second-order form:

- No need for a staggered grid since the operator Δ is elliptic.
- One can solve for \mathbf{E} alone.



Modified Equation time stepping

Taylor series in time:

$$\frac{u(t + \Delta t) - 2u(t) + u(t - \Delta t)}{\Delta t^2} = u_{tt} + \frac{\Delta t^2}{12} u_{tttt} + O(\Delta t^4)$$

For the wave equation

$$u_{tt} = \Delta u$$

a fourth-order scheme in space and time is

$$\frac{U_i^{n+1} - 2U_i^n + U_i^{n-1}}{\Delta t^2} = \Delta_{4h} U_i^n + \frac{\Delta t^2}{12} (\Delta^2)_{2h} U_i^n$$

This scheme is very efficient (especially on Cartesian grids) and allows a large (cfl=1) time step.



Centered numerical boundary conditions for high-order approximations

Vector wave equation on a square

$$\mathbf{E}_{tt} = \mathbf{E}_{xx} + \mathbf{E}_{yy} \quad \mathbf{x} \in \Omega = [0, 1]^2$$

PEC (perfect electrical conductor) boundary at $x = 0$:

$$\begin{aligned} E^y(0, y, t) &= 0 && \text{(from } \mathbf{n} \times \mathbf{E} = 0), \\ \partial_x E^x(0, y, t) &= 0 && \text{(from } \nabla \cdot \mathbf{E} = 0). \end{aligned}$$

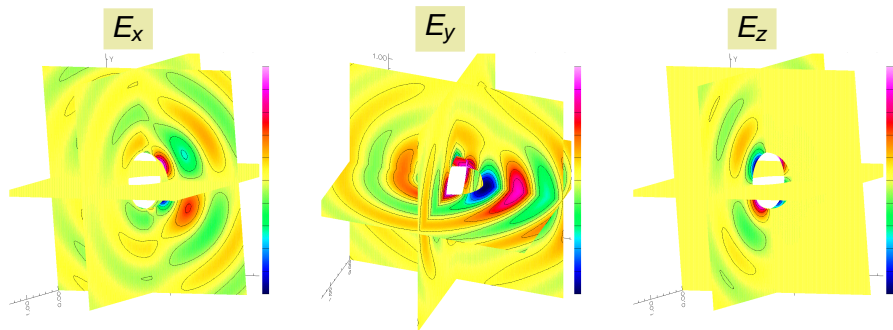
Taking time derivatives of the above and using the equations:

$$\begin{aligned} \partial_x^{2m} E^y(0, y, t) &= 0 \quad m = 0, 1, 2, 3, \dots \\ \partial_x^{2m+1} E^x(0, y, t) &= 0 \quad m = 0, 1, 2, 3, \dots \end{aligned}$$

These *centered* conditions are used on the boundary instead of one-sided approximations.



Scattering of a plane wave by a sphere



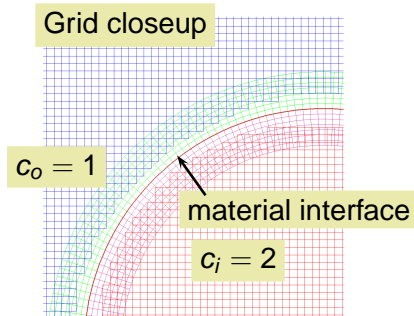
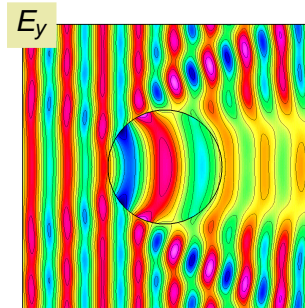
| grid | N | $\ e^{E_x}\ _\infty$ | $\ e^{E_y}\ _\infty$ | $\ e^{E_z}\ _\infty$ | $\ \nabla \cdot \mathbf{E}\ _\infty$ |
|------|-----|----------------------|----------------------|----------------------|--------------------------------------|
| sib1 | 40 | $1.1e-2$ | $7.9e-3$ | $5.6e-3$ | $4.0e-3$ |
| sib2 | 80 | $8.1e-4$ | $5.6e-4$ | $4.0e-4$ | $4.2e-4$ |
| sib4 | 160 | $5.4e-5$ | $3.7e-5$ | $2.7e-5$ | $5.4e-5$ |
| rate | | 3.84 | 3.87 | 3.86 | 3.10 |

Maximum errors at $t = 3$.

The finest grid has 6.5 million grid points.



Scattering of a plane wave by a dielectric cylinder

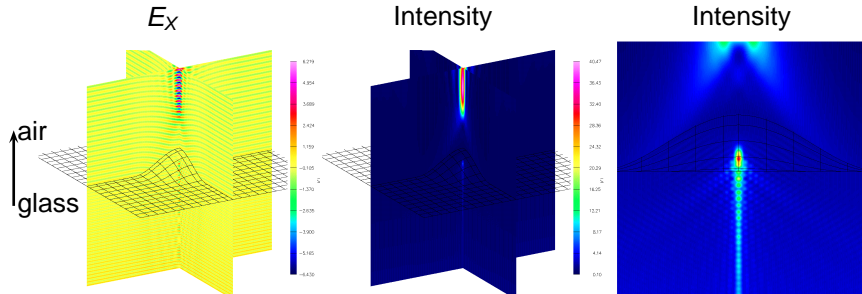


| grid | $\ e^{E_x}\ _\infty$ | $\ e^{E_y}\ _\infty$ | $\ e^{H_z}\ _\infty$ | δ_E |
|-----------------|----------------------|----------------------|----------------------|------------|
| \mathcal{G}_1 | $1.4e-1$ | $2.9e-1$ | $3.0e-1$ | $6.7e-2$ |
| \mathcal{G}_2 | $1.0e-2$ | $2.1e-2$ | $2.2e-2$ | $4.5e-3$ |
| \mathcal{G}_4 | $6.8e-4$ | $1.4e-3$ | $1.4e-3$ | $2.9e-4$ |
| rate σ | 3.86 | 3.87 | 3.88 | 3.92 |

Known solution as a Mie series. Maximum errors at $t = 1$.



Scattering by a 3d material interface

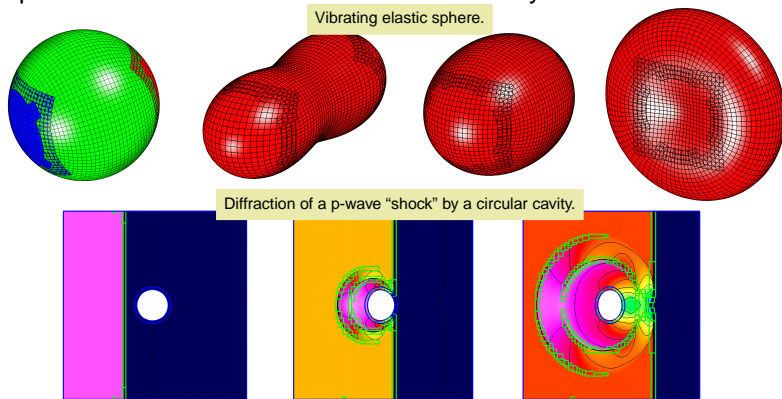


- Uses newly developed 4th-order accurate 3D material interface approximations.
- Scattering of a plane wave by an interface with a bump, glass-to-air.
- 1 billion grid points, 32 nodes (8 processors per node) of a Linux cluster.



Cgsm: a solid-mechanics solver (in Overture.v24).

- linear elasticity on overlapping grids, with adaptive mesh refinement,
- conservative finite difference scheme for the second-order system,
- upwind Godunov scheme for the first-order-system.

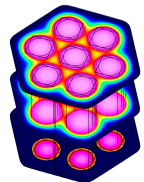
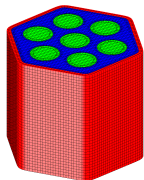


- D. Appelö, J.W. Banks, WDH, D.W. Schwendeman, *Numerical Methods for Solid Mechanics on Overlapping Grids: Linear Elasticity*, LLNL-JRNL-422223, submitted.



Cgmp: a multi-domain multi-physics solver.

Conjugate heat transfer: coupling incompressible flow to heat conduction in solids.

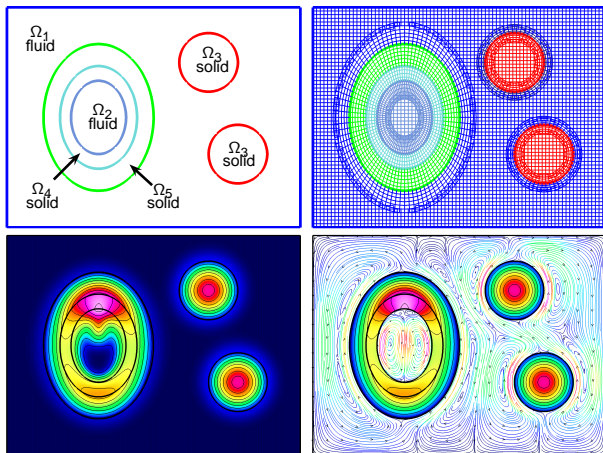


- overlapping grids for each fluid or solid domain,
- a partitioned solution algorithm (separate physics solvers in each sub-domain),
- (cgins) incompressible Navier-Stokes equations (with Boussinesq approximation) for fluid domains,
- (cgad) heat equation for solid domains,
- a key issue is interface coupling.

- WDH., K. K. Chand, *A Composite Grid Solver for Conjugate Heat Transfer in Fluid-Structure Systems*, J. Comput. Phys, 2009.



The multi-domain composite grid approach for CHT



Each fluid or solid sub-domain is covered by an overlapping grid.
Fluid sub-domains : cgins. Solid sub-domains: cgad.
Coupled problem: cgmp.



Deforming composite grids for FSI

Goal: Couple overlapping grid techniques for modeling fluids and gases (using moving grids and AMR) with linear and non-linear solid mechanics codes.

Approach:

- Fluids: Overlapping grid fluid-mechanics solver.
- Solids : unstructured grid or overlapping-grid solid-mechanics solver.
- Boundary fitted deforming grids are used at the fluid-solid interfaces.

Strengths of the approach:

- Maintains high quality grids for large deformations and displacements.
- Uses efficient structured grid methods optimized for Cartesian grids.

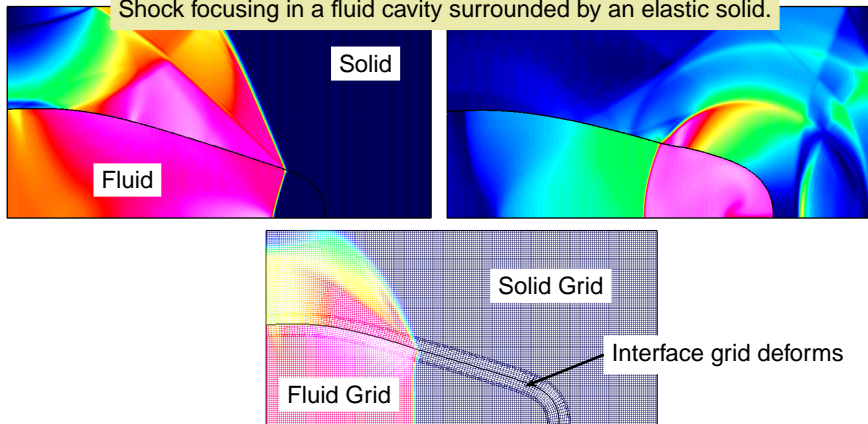
Current status:

- Solve Euler equations in the fluid domains on moving grids.
- Solve equations of linear elasticity in the solid domains.
- Fluid grids at the interface deform over time.



Deforming composite grids for FSI

Shock focusing in a fluid cavity surrounded by an elastic solid.



- Solving the Euler equations in the fluid, linear elasticity in the solid.

The figures show results from preliminary work to model an experiment by Veronica Eliasson.



Conclusions

- Overlapping grids have been used to solve a wide class of problems.
- Smooth boundary fitted grids for accuracy.
- Structured grids for efficiency.
- Rapid grid generation for moving geometry.
- Overture is a toolkit for grid generation and solving PDEs.
- The CG set of PDE solvers solve a variety equation in continuum mechanics .

Open problem: automatic grid generation for complex geometry.

